# Equivariant Seidel maps and a flat connection on equivariant symplectic cohomology

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 $(M, \omega)$  is a closed symplectic manifold.  $\sigma: S^1 \times M \to M$  is a Hamiltonian  $S^1$ -action.

The **clutching bundle**  $E(\sigma)$  is a bundle with fibre *M* over the sphere  $S^2$ .



### The quantum Seidel map

$$QS(\sigma): QH^*(M) \to QH^{*+|\sigma|}(M)$$

counts pseudoholomorphic sections of the clutching bundle  $E(\sigma) \rightarrow S^2$ .

**Novikov ring** is  $\Lambda = \mathbb{Z}[q^{H_2(M)}] \ni q^A$  with  $A \in H_2(M)$ . **Quantum cohomology** is  $QH^*(M) = H^*(M) \otimes \Lambda$ .

#### Seidel maps

Quantum product



**Novikov ring** is  $\Lambda = \mathbb{Z}[q^{H_2(M)}] \ni q^A$  with  $A \in H_2(M)$ . **Quantum cohomology** is  $QH^*(M) = H^*(M) \otimes \Lambda$ .

#### Quantum product is

$$x_1^+ * x_2^+ = \sum_{\substack{A \in H_2(M) \\ x^- \in H^*(M)}} \# \left( \bigvee_{\substack{x_1^+ \cdots x_2^+}}^{x^-} \right) q^A x^-$$

Theorem (Seidel, '97)  
We have 
$$QS(\sigma)(\alpha * x^+) = \alpha * QS(\sigma)(x^+)$$
 for  $\alpha \in QH^*(M)$ .

#### Seidel maps Proof of $QS(\sigma)(\alpha * x^+) = \alpha * QS(\sigma)(x^+)$





#### Seidel maps Proof of $QS(\sigma)(\alpha * x^+) = \alpha * QS(\sigma)(x^+)$







We showed  $QS(\sigma)(\alpha * x^+) = \alpha * QS(\sigma)(x^+)$ . Therefore, we have  $QS(\sigma)(x) = x * QS(\sigma)(1)$ . This yields the **Seidel representation** 

$$\begin{split} \left\{ \begin{aligned} & \text{Hamiltonian } S^1\text{-actions} \right\} \to QH^*(M)^{\times} \\ & \sigma \mapsto Q\mathcal{S}(\sigma)(1). \end{aligned} \tag{1}$$

## Equivariant Seidel maps

Equivariant cohomology





#### Equivariant Seidel maps Definition



In the Borel construction, we have a fibre bundle

 $\frac{S^{\infty}}{\varsigma^1}$ 



S<sup>1</sup>-equivariant quantum Seidel map is

$$Q\mathcal{S}_{S^{1}}(\sigma): QH_{S^{1}}^{*}(M, \sigma) \to QH_{S^{1}}^{*+|\sigma|}(M, \mathrm{Id})$$
$$Q\mathcal{S}_{S^{1}}(\sigma)(\varepsilon^{+}, x^{+}) = \sum_{\substack{A \in H_{2}(M) \\ (\varepsilon^{-}, x^{-}) \in H_{S^{1}}^{*}(M, \mathrm{Id})} \#(\bigcup_{s \in \mathbb{R}^{*}}^{t} e_{s \in \mathbb{R}^{*}}^{t}) q^{A}(\varepsilon^{-}, x^{-})$$



# Theorem (Intertwining relation, TL-J) We have

$$QS_{S^1}(\sigma)(x*\alpha^+) - QS_{S^1}(\sigma)(x)*\alpha^- = \mathbf{u} \ WQS_{S^1}(\sigma,\alpha)(x) \quad (2)$$

for any class  $\alpha \in H^2_{S^1}(E(\sigma))$ . Here,  $\alpha^{\pm}$  are the restrictions of  $\alpha$  to the fibres above the poles,  $\mathbf{u} \in H^2(\mathbb{CP}^{\infty})$  is the generator of  $H^*(\mathbb{CP}^{\infty})$ , and  $WQS_{S^1}$  is a weighted version of the  $S^1$ -equivariant quantum Seidel map.

We proved this for all  $\alpha \in H^*_{S^1}(E(\sigma))$  but we'll only present  $|\alpha| = 2$  for simplicity.

#### Equivariant Seidel maps

 $\mathsf{Proof of} \ \mathcal{QS}_{\mathsf{S}1}(\sigma)(\mathsf{x}\ast\alpha^+) - \mathcal{QS}_{\mathsf{S}1}(\sigma)(\mathsf{x})\ast\alpha^- = \mathbf{u} \ \mathcal{WQS}_{\mathsf{S}1}(\sigma,\alpha)(\mathsf{x})$ 







To parameterise the line of longitude, we would need an  $S^1$ -equivariant map  $S^\infty \to S^1$ .

But none exists. Let  $w \in S^{\infty}$ . The composition  $S^1 \cdot w \hookrightarrow S^{\infty} \to S^1$  is an isomorphism. But  $S^{\infty}$  is contractible, so  $\pi_1(S^{\infty}) = 0$ .

**Key insight:** it is sufficient to define  $S^{\infty} \to S^1$  on a *generic* subset  $W \subset S^{\infty}$ .

#### Equivariant Seidel maps

Proof of  $QS_{\varsigma1}(\sigma)(x * \alpha^+) - QS_{\varsigma1}(\sigma)(x) * \alpha^- = \mathbf{u} WQS_{\varsigma1}(\sigma, \alpha)(x)$ 







Maulik and Okounkov defined equivariant Seidel maps in 2013. They also proved the intertwining relation for  $\alpha \in H^2_{S^1}(E(\sigma))$ . Iritani gave a similar construction in a different setting in 2017.

Their definitions use *virtual fundamental classes* to count sections. Their proofs of the intertwining relation use *virtual localisation*.

We are interested in  $S^1$ -equivariant Floer theory, which does not have the above machinery. We redefined  $QS_{S^1}(\sigma)$  using a Morse Borel construction. We reproved the intertwining relation with a new Morse homotopy proof using a 1-dimensional moduli space argument.



The Floer Seidel map  $FS(\sigma): FH^*(M; H) \to FH^{*+|\sigma|}(M; \sigma^*H)$ maps the Hamiltonian orbit  $x: S^1 \to M$  to the orbit  $(\sigma^*x)(t) = \sigma_t^{-1}(x(t)), \quad t \in S^1.$ 

 $FS(\sigma)$  is an isomorphism of cochain complexes.

**Compact:**  $QH^*(M) \cong FH^*(M; H)$  for all Hamiltonians *H*. The Floer Seidel map and the quantum Seidel map agree.

**Non-compact:** (with convexity assumption, for example  $\mathcal{O}_{\mathbb{P}^1}(-1)$ )  $FH^*(M, \lambda; H)$  depends on the slope  $\lambda$  ( $H = \lambda R + \text{const. at infinity}$ ) **Symplectic cohomology** is  $SH^*(M) = \varinjlim_{K} FH^*(M, \lambda)$  as  $\lambda \to \infty$ . For *linear*  $\sigma$  of slope  $\kappa$ , the Floer Seidel map is  $FS(\sigma) : FH^*(M, \lambda; H) \to FH^{*+|\sigma|}(M, \lambda - \kappa; \sigma^*H)$ .

The quantum Seidel map is only defined for linear  $\sigma$  with  $\kappa \geq 0$ .



The loop space  $\mathcal{L}M = \{x: S^1 \to M\}$  has an  $S^1$ -action given by

$$(\theta \cdot x)(t) = x(t-\theta) \qquad \theta \in S^1.$$
 (3)

## Definition (Equivariant Floer cohomology)

We combine Morse theory on  $S^{\infty}$  with Floer theory on M. The Hamiltonian  $H: S^{\infty} \times S^1 \times M \to \mathbb{R}$  now depends on  $S^{\infty}$ . The **equivariant Floer cochain complex** is generated over  $\Lambda$  by  $[\varepsilon, x]$ , where  $\varepsilon \in S^{\infty}$  is critical and x is a Hamiltonian orbit of  $H_{\varepsilon}$ . The differential counts flowlines in  $S^{\infty}$ paired with Floer cylinders in M.

$$FH^*_{S^1}(M,\lambda;H)$$
 is a  $\Lambda\otimes\mathbb{Z}[\mathbf{u}]$ -module.



There is a map  $QH^*(M) \otimes FH^*(M, \lambda) \to FH^*(M, \lambda)$  which counts



An equivariant version of this map would use a map  $S^{\infty} \to S^1$  defined on a *generic* subset  $W \subset S^{\infty}$ . Therefore it would not be a chain map.

## Equivariant Floer cohomology

Differentiation



**Novikov ring** is  $\Lambda = \mathbb{Z}[q^{H_2(M)}] \ni q^A$  with  $A \in H_2(M)$ . For  $\alpha \in H^2(M)$ , define

$$\frac{d}{d\alpha}(q^A) = \alpha(A) \ q^A. \tag{4}$$

We can pick a  $\Lambda$ -basis for the equivariant Floer cochain complex and apply  $\frac{d}{d\alpha}$ . This is not a chain map either.





On  $QH^*_{S^1}(M, \mathrm{Id}) = \Lambda \otimes \mathbb{Z}[\mathbf{u}] \otimes H^*(M)$ , the **Dubrovin connection** is

$$\nabla_{\alpha}(q^{A}x) = \mathbf{u}\frac{d}{d\alpha}(q^{A})x + \alpha * q^{A}x.$$
 (5)

Theorem (Connection, TL-J) On  $FH^*_{S^1}(M, \lambda; H)$ , for  $\alpha \in H^2(M)$  the map

$$\nabla_{\alpha} = \mathbf{u} \frac{d}{d\alpha} + (\alpha * \cdot) - w_{\alpha} \tag{6}$$

is a chain map on the equivariant Floer cochain complex.

Seidel proved special case  $\alpha = [\omega]$  in 2016.

Connection and Seidel maps

Theorem (Flatness, TL-J) On  $FH^*_{S^1}(M, \lambda; H)$ , for any  $\alpha, \beta \in H^2(M)$ , we have  $\nabla_{\alpha} \circ \nabla_{\beta} = \nabla_{\beta} \circ \nabla_{\alpha}.$ 

The Dubrovin connection is  $\nabla_{\alpha} = \mathbf{u} \frac{d}{d\alpha} + \alpha *$ . A calculation shows the intertwining relation

 $QS_{S^1}(\sigma)(x * \alpha^+) - QS_{S^1}(\sigma)(x) * \alpha^- = \mathbf{u} \ WQS_{S^1}(\sigma, \alpha)(x)$ is equivalent to

$$Q\mathcal{S}_{S^1}(\sigma)(\nabla_{\alpha}(x)) - \nabla_{\alpha}(Q\mathcal{S}_{S^1}(\sigma)(x)) = 0.$$

Theorem (Connection and Floer Seidel map, TL-J) On  $FH^*_{S^1}(M, \lambda; H)$ , for any  $\alpha \in H^2(M)$  and any linear  $\sigma$ , we have  $\nabla_{\alpha} \circ FS_{S^1}(\sigma) = FS_{S^1}(\sigma) \circ \nabla_{\alpha}.$ 



Torus action



Now let T be a torus acting on M.  $S^1$ -actions correspond to **cocharacters** of T, which are group maps  $\sigma : S^1 \to T$ . Let  $\widehat{T} = S_0^1 \times T$ . Constructions of  $QH^*_{\widehat{T}}(M)$ ,  $FH^*_{\widehat{T}}(M, \lambda)$ ,  $SH^*_{\widehat{T}}(M)$ . We get  $\nabla_{\alpha}$ ,  $QS_{\widehat{T}}(\sigma)$ ,  $FS_{\widehat{T}}(\sigma)$  too. But now we can undo the change of  $\widehat{T}$ -action, to get endomorphisms

$$S_{\sigma} : QH^*_{\widehat{T}}(M) \to QH^{*+|\sigma|}_{\widehat{T}}(M)$$
$$S_{\sigma} : SH^*_{\widehat{T}}(M) \to SH^{*+|\sigma|}_{\widehat{T}}(M)$$
(7)

#### called shift operators.

#### Shift operators Example: P<sup>2</sup>



The torus  $T^2$  acts on  $\mathbb{P}^2$  (on middle and last coordinate). We have  $H^*(BT) = \mathbb{Z}[t_1, t_2]$  and  $\Lambda = \mathbb{Z}[q]$  with |q| = 6.

$$QH^*_{\widehat{T}}(\mathbb{P}^2) = \frac{\Lambda \otimes \mathbb{Z}[x_0, x_1, x_2, \mathbf{u}]}{x_0 x_1 x_2 - q}$$
(8)

We calculate  $\nabla_x = \mathbf{u}(q\frac{d}{dq}) + x_0$ . Let  $\sigma$  correspond to rotation of middle coordinate. We have:

$$S_{\sigma}(1) = x_{1} \qquad S_{\sigma}(t_{1}y) = (t_{1} + \mathbf{u})S_{\sigma}(y)$$

$$S_{\sigma}(x_{0}) = x_{1}x_{0} \qquad S_{\sigma}(t_{2}y) = t_{2}S_{\sigma}(y)$$

$$S_{\sigma}(x_{1}) = x_{1}(x_{1} - \mathbf{u})$$

$$S_{\sigma}(x_{2}) = x_{1}x_{2}$$
(9)



### We had the Seidel representation

$$\begin{cases} \mathsf{Hamiltonian} \ S^{1}\text{-actions} \end{cases} \to QH^{*}(M)^{\times} \\ \sigma \mapsto QS(\sigma)(1). \end{cases}$$
(10)

We also have 
$$\mathbb{S}_{\sigma}\mathbb{S}_{\sigma'} = \mathbb{S}_{\sigma+\sigma'}$$
.  
This yields

$$\mathbb{S} : \operatorname{Cochar}^+(T) \to \operatorname{End}_{\Lambda[\mathbf{u}]}(QH^*_{\widehat{T}}(M))$$
  
 $\mathbb{S} : \operatorname{Cochar}(T) \to \operatorname{Aut}_{\Lambda[\mathbf{u}]}(SH^*_{\widehat{T}}(M)).$ 



We have expanded the algebraic structures on  $SH^*_{\widehat{\tau}}(M)$  with

- ▶ a flat connection  $\nabla_{\alpha}$ ,
- shift operators  $\mathbb{S}_{\sigma}$ .

They are compatible, computable in examples and capture geometric information.

Thanks for your attention.



Definition (Borel construction of  $S^1$ -equivariant cohomology) X is topological space,  $\rho$  is  $S^1$ -action on X. Take  $S^\infty$ : it's a contractible space with a free  $S^1$ -action.

**Borel quotient** is  $S^{\infty} \times X / \sim$ , where  $\sim$  is  $(\theta \cdot w, x) \sim (w, \rho_{\theta}(x))$ . It's denoted  $S^{\infty} \times_{S^1} X$ .

 $S^1$ -equivariant cohomology is  $H^*_{S^1}(X, \rho) = H^*(S^{\infty} \times_{S^1} X)$ .

The projection map  $S^{\infty} \times_{S^1} X \to S^{\infty}/S^1 = \mathbb{CP}^{\infty}$  induces a map  $H^*(\mathbb{CP}^{\infty}) = \mathbb{Z}[\mathbf{u}] \to H^*_{S^1}(X, \rho).$ 



The Floer cohomology  $FH^*(M)$  is inspired by Morse cohomology on the loop space  $\mathcal{L}M = \{x : S^1 \to M\}$ .

Take a function  $H: S^1 \times M \to \mathbb{R}$ , called a **Hamiltonian function**. Define the **Hamiltonian vector field** by  $\omega(\cdot, X_t) = dH_t$ . The **Hamiltonian orbits** are the curves  $x: S^1 \to M$  that follow  $X_t$ . The **Floer cochain complex** is freely generated over  $\Lambda$  by the Hamiltonian orbits.

A Floer cylinder is a cylinder  $u : \mathbb{R} \times S^1 \to M$ which satisfies a Floer equation. The Floer differential counts Floer cylinders between the Hamiltonian orbits.





# Definition (Convex symplectic manifold)

A convex symplectic manifold is the union of a compact symplectic manifold and the symplectic manifold  $([1,\infty) \times \Sigma, d(R\alpha))$ , where  $(\Sigma, \alpha)$  is a closed contact manifold. Example  $\mathcal{O}_{\mathbb{P}^1}(-1)$ , where  $\Sigma = S^3$  is the sphere bundle.

Floer cohomology depends on the **slope**  $\lambda$ , where  $H = \lambda R$  + constant at infinity.

Symplectic cohomology  $SH^*(M)$  is the limit of  $FH^*(M, \lambda)$  as  $\lambda \to \infty$ .